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THERMODYNAMIC PROPERTIES OF CESIUM UP TO 1500° K

by Sheldon Heimel Lewis Research Center Cleveland, Ohio

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NATIONAL AERONAUTICS AND SPACE ADMINISTRATION

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THERMODYNAMIC PROPERTIES OF CESIUM UP TO 1500° K

by Sheldon Heimel

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SUMMARY

Consistent tables of thermodynamic properties of elemental cesium were compiled for both the pure species and the equilibrium vapor on the saturation line using selected values of 10 500 calories per mole for the heat of dissociation of dimer and -18 920 calories per mole for the heat of condensation of monomer. The equilibrium vapor properties and properties of condensed cesium (Cs) are given to 1500° K, while properties of gaseous monomer and dimer are given to 2500° K.

The table of mixture properties on the saturation line gives the following thermodynamic functions: temperature T, total vapor pressure P, weight fraction of dimer in vapor x_2 , molecular weight M_g , specific volume v, density ρ , enthalpy of condensed phase h_c , enthalpy of vaporization Δh_{vap} , enthalpy of vapor h_g , entropy of condensed phase s_c , entropy of vaporization Δs_{vap} , entropy of vapor s_g , and heat capacities (equilibrium and frozen) $(c_p)_{eq}$ and $(c_p)_{fr}$. Assigned enthalpies are relative to enthalpy of the crystal at 298.15° K, which is taken to be zero. The tables of properties for the pure species give the following thermodynamic functions for the standard state: heat capacity at constant pressure C_p^0 , sensible enthalpy $H_T^0 - H_O^0$, entropy S_T^0 , sensible free energy $F_T^0 - H_O^0$, the sum of sensible enthalpy at T^0 K and chemical energy at 0^0 K H_T^0 , and values of enthalpy changes and logarithms of the equilibrium constants. The latter two functions are given for the reactions of formation of Cs(c), $Cs_1(g)$, and $Cs_2(g)$ from the assigned reference state Cs(c) ((ΔH_T^0) and log_{10} K_f) and from $Cs_1(g)$ (ΔH_T^0 and log_{10} K).

Thermodynamic functions for the gases were generated from atomic and molecular data, whereas the functions of the condensed phase were based on selected experimental data that were smoothed and made self-consistent.

INTRODUCTION

Current work at the Lewis Research Center indicates the need for reliable vapor pressure data as well as other thermodynamic properties for cesium in the following areas:

- (1) Rankine cycle engine. The thermodynamic properties of cesium are needed in order to make an analysis of a Rankine cycle system and its components.
- (2) Thermionic plasma diode. Improved performance requires knowledge of the concentration of cesium in the gas phase and the rate of evaporation from a liquid surface. Such data can be derived with the help of an equation for pressure as a function of saturation temperature.

Accordingly, to meet these needs, the following tables have been generated:

- (1) Boiling points corresponding to pressures of from 10^{-8} to 10 atmospheres
- (2) Properties on the saturation line from $100^{\rm O}$ to $1500^{\rm O}$ K: temperature T, $^{\rm O}$ K; total vapor pressure P, atmospheres; partial pressure of monomer $\rm p_1$, atmospheres; partial pressure of dimer $\rm p_2$, atmospheres; weight fraction of dimer $\rm x_2$, gram per gram mixture; molecular weight of vapor $\rm M_g$, gram per mole; specific volume of vapor v, cubic centimeter per gram; density of vapor ρ , gram per 1000 cubic centimeters; enthalpy of condensed phase $\rm h_c$, calories per gram; enthalpy of vaporization $\Delta \rm h_{vap}$, calories per gram; enthalpy of vapor $\rm h_g$, calories per gram; entropy of condensed phase $\rm s_c$, calories per gram per $\rm ^{\rm O}$ K; entropy of vaporization $\Delta \rm s_{vap}$, calories per gram per $\rm ^{\rm O}$ K; entropy of vapor $\rm s_g$, calories per gram per $\rm ^{\rm O}$ K; equilibrium heat capacity of vapor ($\rm c_p$) $\rm _{eq}$, calories per gram per $\rm ^{\rm O}$ K; and frozen heat capacity of vapor ($\rm c_p$) $\rm _{eq}$, calories per gram per $\rm ^{\rm O}$ K; and frozen heat capacity of vapor ($\rm c_p$) $\rm _{eq}$, calories per gram per $\rm ^{\rm O}$ K
- (3) Thermodynamic functions for each pure species: temperature T, O K; heat capacity at constant pressure C_p^o , calories per mole per O K; sensible enthalpy $H_T^o H_O^o$, calories per mole; entropy S_T^o , calories per mole per O K; sensible free energy $F_T^o H_O^o$, calories per mole; sum of sensible enthalpy at T^O K and chemical energy at O K H_T^o , calories per mole; and values of enthalpy changes and logarithms of the equilibrium constants. The latter two functions are each given for two different reactions of formation. One set of functions, $(\Delta H_T^o)_f$ and $\log_{10} K_f$ is for the reaction of formation of $Cs_1(g)$ and $Cs_2(g)$ from the assigned reference state Cs(c). The other set of functions, ΔH_T^o and $\log_{10} K$, is for the formation of Cs(c) and $Cs_2(g)$ from $Cs_1(g)$. The symbols g and c refer to the gas and condensed phases (crystal and liquid), respectively. The thermodynamic functions for the gases were generated from atomic and molecular data, whereas the functions of the condensed phase were

based on selected experimental data that were smoothed and made self-consistent. The functions for the condensed phase are given up to 1500° K; those for the monomeric and dimeric gases are given up to 2500° K.

When vapor pressure data are available, they may be combined with sensible free energy data by use of the third law of thermodynamics to calculate an implied value of $(\Delta H_O^0)_c$ (heat of condensation of monomer at 0^0 K) for each available vapor pressure point. If the vapor pressure data are completely consistent with the thermodynamic data, the calculated values of $(\Delta H_O^0)_c$ will be constant. The usual result, however, is that there is variation in the calculated values of $(\Delta H_O^0)_c$. Depending on the extent and nature of this variation, some average $(\Delta H_O^0)_c$ may be selected. This selected value of $(\Delta H_O^0)_c$ together with the sensible free energy data may now be used to calculate vapor pressures. These calculated vapor pressures and the experimental vapor pressures will, of course, not agree. At this point either the experimental vapor pressures must be abandoned or the theoretical free energies must be modified.

In this report a similar procedure to that just outlined will be followed. There are several differences in the analysis, namely,

- (1) The experimental vapor pressures are first smoothed.
- (2) The gas is assumed to be a mixture of monomer and dimer. Therefore, the intermediate step of calculating equilibrium compositions is required.
- (3) The procedure for selecting a value of $(\Delta H_O^O)_c$ is to find that value of $(\Delta H_O^O)_c$ which gives a minimum deviation of calculated and smoothed experimental vapor pressures over the entire range.
- (4) The selected $(\Delta H_O^0)_c$ and thermodynamic data are used to calculate a recommended set of vapor pressures for the crystal and the liquid.

PRESENTATION OF TABLES

Table I contains the smoothed values of C_p^O for crystalline cesium up to 100^O K together with consistent values of H_T^O - H_O^O , S_T^O , and -(F_T^O - H_O^O). Generation of these data will be discussed in SELECTION OF INITIAL DATA.

Tables II to IV (pp. 5-7) give the thermodynamic properties of Cs (crystal, liquid) up to $1500^{\rm O}$ K; ${\rm Cs_1(g)}$ up to $2500^{\rm O}$ K; and ${\rm Cs_2(g)}$ up to $2500^{\rm O}$ K, respectively. The properties presented are ${\rm C_p^O}$, ${\rm H_T^O}$ - ${\rm H_O^O}$, ${\rm S_T^O}$, -(F $_{\rm T}^{\rm O}$ - H $_{\rm O}^{\rm O}$), ${\rm H_T^O}$, ($\Delta{\rm H_T^O}$), $\log_{10}{\rm K_f}$, $\Delta{\rm H_T^O}$, and $\log_{10}{\rm K}$. In table II, data for the crystal extend up to the melting point, 301.8 $^{\rm O}$ K. Data for the liquid are tabulated at $100^{\rm O}$ intervals from the melting point to $1500^{\rm O}$ K. In

TABLE I. - THERMODYNAMIC PROPERTIES

OF CESIUM (CRYSTAL) UP TO 100⁰ K

_				
T,	C _p ,	$H_{\mathbf{T}}^{\mathbf{O}} - H_{\mathbf{O}}^{\mathbf{O}}$	$\mathbf{S_{T}^{O}},$	$-(\mathbf{F}_{\mathbf{T}}^{\mathbf{O}} - \mathbf{H}_{\mathbf{O}}^{\mathbf{O}}),$
°К	cal/(mole)(^O K)	cal/mole	cal/(mole)(^O K)	cal/mole
5	0.746	1. 2	0, 3342	0.5
10	2.379	9.0	1, 3539	4.5
15	3.719	23.9	2,5835	14.9
20	4.680	44.5	3.7993	31.5
25	5.110	69. 1	4,8935	53.3
30	5.385	95.3	5.8513	80.2
35	5.558	122.7	6.6955	111.6
40	5.668	150.8	7.4454	147.0
45	5.741	179.4	8. 1174	185.9
50	5.795	208. 2	8.7252	228. 1
55	5.842	237.3	9. 2798	273.1
60	5.888	266.6	9. 7901	320.8
65	5.933	296. 2	10. 2631	370.9
70	5.977	325.9	10. 7044	423.4
75	6.016	355.9	11. 1182	477.9
80	6.049	386.1	11.5076	534.5
85	6.075	416.4	11.8751	593.0
90	6.097	446.8	12. 2230	653.2
95	6.122	477.4	12. 5533	715.2
100	6.160	508.1	12.8682	778.7

tables III (Cs₁) and IV (Cs₂), data are tabulated at 100⁰ intervals from 0⁰ to 2500⁰ K for the ideal gases. Table V (p. 8) contains the molecular constants used to calculate thermodynamic functions for dimeric cesium.

Table VI (p. 8) contains recommended boiling points of liquid cesium calculated from the final set of thermodynamic properties. These boiling points correspond to pressures from 10⁻⁸ to 10 atmospheres at every power of 10. For comparison, boiling points obtained from equation (33) are also given.

Table VII (p. 9) gives the following thermodynamic properties of the equilibrium gaseous mixture on the saturation line: T, P, p_1 , p_2 , x_2 , M_g , v, ρ , h_c , Δh_{vap} , h_g , s_c , Δs_{vap} , s_g , $(c_p)_{eq}$, and $(c_p)_{fr}$. The values of P, p_1 , p_2 , v, and ρ are given in floating-point notation, where the decimal number is to be multiplied by

10 raised to the power of the sign and the two digits following the letter E; for example, 0.100000E-03 is $0.100000\times10^{-3} = 0.0001$.

Assigned Reference State

The assigned reference state is Cs (crystal, liquid). The crystal is the reference state up to the melting point, and the liquid phase is the reference state above the melting point.

Assigned Enthalpy Values Ho

For some applications (ref. 1), it is convenient to combine sensible enthalpy ${\rm H}_{\rm T}^{\rm O}$ - ${\rm H}_{\rm O}^{\rm O}$ and chemical energy ${\rm H}_{\rm O}^{\rm O}$ into one numerical value, ${\rm H}_{\rm T}^{\rm O}$. The arbitrary base

TABLE II. - THERMODYNAMIC PROPERTIES OF CESIUM (CRYSTAL, LIQUID)

o T,	$C_{\mathbf{p}}^{0}$	$H_{\rm T}^{\rm O}$ - $H_{\rm O}^{\rm O}$, cal/mole	S.T.	$-(F_{\rm T}^{\rm O}-H_{\rm O}^{\rm O}),$	$_{ m H}^{ m O}_{ m T}$	Formation from assigned reference state	om assigned se state	Formation from gaseous atoms	on from atoms
	(mole)(^O K)		(Mole)(OK)			$(\Delta H_{T}^{O})_{\mathbf{f}},$ cal/mole	$^{\log_{10}~\mathrm{K_f}}$	$\Delta H_{\mathrm{T}}^{\mathrm{O}},$ cal/mole	$\log_{10} m K$
0		G	o	0	-1307.2	C	1	r*18851-	†
100	6.160	500.1	12.8602	176.1	1.9951-	9		-1.5905-7	34.1542
200	6.560	1144.1	17.2607	2308.0	-063.1	c			15.54%
258.15	6.953	1807.2	7856.61	4141.8	○• 0	0		-13594.7	6,600,60 10,000,000
300	095*9	1520.1	19.4962	4178.4	12.3	•	-	-14500	4.73.5
a 301.8	6.507	1832.6	20.0379	4214.8	25.4	,,	-	-18586.3	4.65.4
a361.8	7.455	2352.7	21.7012	4214.8	545.5	n	•	-13064.7	4×5×4
400	7.455	3€84•8	23.4513	6459.7	1277.6	0	·	-17022.4	5.4664
200	7.455	3430•3	25.5248	8932.1	2023.1	2		-17573-4	3. 5315
600	7.455	4575° 8	26.8840	11554.0	2708.6	0	_	-17375-1	2.2504
700	7.455	5321.3	28.0332	14301.3	3514.1	c	-	-17076-4	1.35.45
800	7.455	6366.0	25.0287	17156.1	4259.6	Ċ.	<i>~</i> -	-16827.7	1611°C
006		6812.3	7906.67	20103.3	5004.1	<u>.</u>	1.	-16575.	0.1956
b937.71	7.455	7093.4	30.2127	21237.4	5286.2	9	es.	-16485.2	7.1341
0001		7551.8	36.6922	23134.4	2750.6	0	a	-15330.3	-5.2341
1100	7.455	8303.3	31.4027	26239.7	6436.1	е	** * **	-15081.7	1765-0-
1200	7.455	9048.8	32.0514	29412.9	7241.6	Ċ	63	-15833.1	4CFL-C-
1300	7.455	9794.3	32.6481	32648.3	1987.1	¢	C.	-15584.6	-I.015
1400	7.455	10539.8	33.2006	35941.0	8732.6	ر	<	-15336.4	-1.1.1.2
1500	7.455	11285.3	33.7150	39287.1	9478.1	0)	-15085.7	-1.3545

 $^{^{\}mathbf{a}}\mathbf{Melting}$ point. $^{\mathbf{b}}\mathbf{Normal}$ boiling point to equilibrium mixture.

TABLE III. - THERMODYNAMIC PROPERTIES OF Cs_1 (GAS)

													_	_									_		_			_	_	
Formation from gaseous atoms	$\log_{10} \mathrm{K}$		† - -	c	, بــــــ	<u>~ (</u>	c (5 6		-	٠	7.	· ·	c «	٠,	·	0	٠ ٠	c ′	- (C 1	. ·	Į :	£* 5	. .	= 1	c (- 6	
Formation fron gaseous atoms	$^{\Delta H_{\mathrm{T}}^{\mathrm{O}}}$	cal/mole	C.	Ġ	c	c ·	G-1	<u> </u>		ř.	c	¢.	ς Ι	C: (ς (ç.	O	c :	c (ς. (0.	C		De f	(C I	c (. (. c
om assigned e state	log ₁₀ K _f		1 1 1 1 1	- 37 - 1548	-15.5486	-4.4732	- 3.7392	-4.6534	4064	186.6-	-2.2593	-1.3545	7207-0-	- 1356	- J • G 341	1.2941	0.5251	9L62*0	1.0195	1.1952	1.3545									
Formation from assigned reference state	$(\Delta H_{T}^{O})_{f}$,	cal/mole	18920.0	18908.7	18709.5	18594.0	14590.3	13006.7	17622.4	17573.8	17325.1	17076.4	16027.7	16579.0	16485.2	16330.3	16981.7	15833,1	15584.6	15336.4	15088.7									
$^{ m H}_{ m T}^{ m O}$	car/ more		17112.8	17609.6	14104.4	18594.0	13603.2	19612.2	19130.0	14596.8	20093.7	202507	21087.3	21584.1	217715	22280.4	22577.8	23074.7	23571.7	54:169.0	24556.7	25005.4	25565.3	26067.2	26571.7	27.79.9	5.7837.5	79111117	28636.5	29170.3
$-(F_{\mathrm{T}}^{\mathrm{O}}-H_{\mathrm{O}}^{\mathrm{O}}),$	car, more		0	3154.9	6998.5	11324.4	11102.0	111177.5	15374.4	19772.3	24270.2	78851.4	33503.7	38213.3	6.01004	42938.3	47808.0	52572.9	57579.3	62524.0	07504.2	12511.1	77502.3	82636.5	87738.5	92867.4	98021.8	103200.6	108403.2	113628.8
$ ho_{ m T}^{ m o}$	(mole)(^o K)			36.5168	35,9605	41.9442	41.9749	45.0046	43.4042	44.5128	45.4186	46.1844	46.8478	47.4330	47.0369	47.9564	48.4300	48.5023	45.2601	45.0287	45.9721	50.2939	60.5970	50.8839	51.1566	51.4173	51.5674	51.9086	52-1422	52-3653
$H_{\rm T}^0 - H_{\rm O}^0$	can/ more		ú	8.055	993.0	1481.3	1490.4	5. 6651	1587.3	2484.1	5*2857	3477.7	3974.5	4471.3	4650.7	4908.2	5465.0	5561.9	6458.9	6950.2	7454.0	7952.6	8452.6	8954.4	6459.0	1.1955	10479.8	10998.3	11523.8	12057.5
$c_{\mathrm{p}}^{\mathrm{o}}$	(mole)(^o K)			4.5001	4.5681	4.5601	4.5681	4.9681	4.5081	4.9641	4.5681	4.5682	4.9682	4.5682	4.5682	4.9083	4.9686	4.9654	4.9712	6515.4	4.9814	5155*5	5.0079	5.0307	5.0618	5.1025	5.1541	5.2175	5.2536	5.3832
T, ^o K	-		12	- 3	200	298.15	300	361.8	3 4€	200	909	700	900	006	937.71	1000	1100	1200	1360	1400	1500	1600	1700	1800	1900	2000	2100	2200	2300	2400

 $^{
m a}{
m The}$ reference element is crystal cesium up to 301, 8 $^{
m O}$ K and liquid cesium above 301, 8 $^{
m O}$ K.

TABLE IV. - THERMODYNAMIC PROPERTIES OF Cs_2 (GAS)

T, OK	ļ								
	ည်း	$H_{\rm T}^0 - H_{\rm O}^0$	$ m S_{T}^{o}$	$-(F_{\mathrm{T}}^{0}-H_{\mathrm{O}}^{0}),$	${ m H}_{ m T}^0$	Formation fi	Formation from assigned	Formation from	n from
	[83]	cal/mole	cal	cal/mole	cal/mole	referen	reference state	gaseous atoms	atoms
	(mole)(^O K)		(mole)(OK)			$(\Delta H_{\mathrm{T}}^{\mathrm{O}})_{\mathrm{f}},$	$\log_{10} m K_{ m f}$	$\Delta H_{T}^{O},$	$\log_{10} m K$
						cal/mole		cal/mole	
O		O	1 1	С	23125.6	27340.1	* * * 1	-10500.0	1 1 1
100	8.5436	843.2	57,9961	4,0564	24565.8	27157.1	-52,3203	-10653.3	6986.61
707	9.0457	1743.6	64.2330	11103.0	75463.2	267.45.5	-21.7462	-10743.6	9.3110
298.15	9.1202	2035.4	67.8558	17597.0	76361.0	26361.0	-13-2133	-10827.1	4.4322
300	9.1215	2052.3	67.9162	17722.6	26377.9	26352.1	-15.0041	-10828.6	4.3847
301.8	5.1227	7-8997	2016-10	17844.9	26394.3	25303.3	-12.9797	-10830.0	4.3372
a 40€	5.1853	3567.8	70.5456	24652 • 0	21293.4	24738.2	-8.5271	-10906.7	7.4057
500	9.2572	4490•1	72.6074	31813.6	28215.7	24169.6	-5.8528	-10078.0	1.2102
009	9-325-6	5419.3	74.3012	39161.5	59144.9	73607.7	-4.1114	-11042.4	0.4733
7.00	9.3959	6355.4	75.7441	46565.5	30081.0	23052.8	-2.4967	-11100.0	-0.1578
800	4.4¢74	7298.5	77.0034	54304.2	31024.1	22504.9	£100.7-	-11150.5	6169*0-
006	5.5402	8248.9	16.1227	62061.5	31974.5	21964.3	-1.3321	-11193.7	-0.9410
937.71	9.56el	86098-2	18.5149	65015.0	32334.8	21762.3	-1.1186	-1120k.1	-1.2514
0001	9.6146	950055	75.1317	69925.1	52937.2	21431.0	-0.8051	-11229.6	-1.2132
1100	5.6905	10171.8	80.0516	17864.9	33097.4	20905.3	-0.3343	-11258.1	-1.4365
1200	4.1679	11144.7	50.8981	85932.9	34670.3	20387.2	-0.0424	-11279.0	-1.5232
1300	5.846d	12125.5	61.6830	3.79046	35851.1	6*91861	1067.0	-11292.3	-1.7313
1400	9.9273	13114.2	32.4151	102267.8	36839.8	19374.6	6.4754	-11298.2	-1.9169
1500	10.0053	14111.0	83-1034	110544.1	37836.6	18830.4	0.6745	-11295.9	-2.0345
1,600	10.0528	15116.1	83.7520	118887.2	38841.7			-11289.1	-2.1373
1700	10.1779	16129.6	84.3664	127293.4	39855.2			-11275.4	-2.2730
1800	10.2646	17151.7	84.9506	135759.4	40877.3			-1125/11	-2.3085
1900	10.3527	18182.6	85.5080	144282.6	41908.2			-11235.3	-2.3803
2000	10.4425	19222.3	86.0413	152860.2	42947•9			-11211-8	-2.4449
7100	10.5337	20271.1	86.5530	161490.1	43996.7			-11188.5	-2.5032
2200	10.6266	21329.1	87.0451	170170.2	45054.7			-11167.5	-2.5560
2300	10.7209	22396.5	87.5190	178498.6	46122.1			-11151.0	-2.6942
2400	10.8168	3473.	87.9779	167673.6	47198.9			-11141.6	-7.6483
2500	10.9143	24559.9	88.4214	196493.6	48285.5			-11141.8	-2.6849

 $^{\mathrm{a}}$ The reference element is crystal cesium up to 301.8 $^{\mathrm{o}}$ K and liquid cesium above 301.8 $^{\mathrm{o}}$ K.

TABLE V. - MOLECULAR

CONSTANTS FOR DI-

ATOMIC CESIUM

Molecular weight	265, 82
Symmetry number	2
Electronic state ^a	¹ Σ
Statistical weight	1
$\omega_{\rm e}^{\rm cm^{-1}a}$	41.990
$\omega \times cm^{-1}$	0.08005
$\omega_{\rm o}$ y _o , cm ⁻¹	-0.0001643
B . cm ⁻¹⁵	0.01272
$\alpha_{\rm o}, {\rm cm}^{-1}$	0.000035
D _e , cm-1 ^b	0.467×10 ⁻⁸

^aConstant from ref. 10.

TABLE VI. - BOILING POINTS OF LIQUID CESIUM FROM 10^{-8} TO 10 ATMOSPHERES

log ₁₀ P _{atm}	Boilin	g points, ^O K
	Recommended	From equation (33) ^a
-8	317.8	317.8
-7	345.8	345.8
-6	379, 2	379. 4
-5	420.1	420.3
-4	471. 2	471. 4
-3	537.0	537. 2
-2	625.0	625.0
-1	749. 1	748. 7
0	937. 7	937. 3
1	1264.0	1264.5

 $a_{\log_{10} P_{atm}} = \frac{-4053.30}{T} + 7.04453 - 0.915282 \log_{10} T.$

for assigning values to the enthalpy of Cs(c) was a value of zero at 298.15° K. Since table II gives $(H_{298.15}^O - H_O^O)_{Cs(c)} = 1807.2$ calories per mole, then $(H_O^O)_{Cs(c)} = -1807.2$ calories per mole. From table II, $(\Delta H_O^O)_{Cs(c)} = (H_O^O)_{Cs(c)} - (H_O^O)_{Cs_1(g)} = -18920.0$ calories per mole. Therefore, $(H_O^O)_{Cs_1(g)} = -1807.2 - (-18920.0) = -19920.0$ Cos₁(g) = -19920.0 calories per mole. From table IV, $(\Delta H_O^O)_{Cs_1(g)} = -D_O^O = (H_O^O)_{Cs_2(g)} - 2(H_O^O)_{Cs_1(g)} = -19920.0$ Cos₁(g) = -19920.0 calories per mole. Therefore,

$$(H_O^0)_{Cs_2(g)} = -10\ 500.0 + 2(17\ 112.8) = 23\ 725.6$$
 cal/mole

Heats of Formation

Two sets of values for heats of formation are given in tables II to IV. The first set (col. 7) is for the formation of the given species from Cs(c), (ΔH_T^O) ; the second set (col. 9) is for the formation from $Cs_1(g)$, (ΔH_T^O) . For $Cs_2(g)$ at 298.15° K, for example,

$$(\Delta H_{298.15}^{O})_{f, 2} = (H_{298.15}^{O})_{Cs_{2}(g)} - 2(H_{298.15}^{O})_{Cs(c)}$$
 (1)

bConstant from ref. 11.

TABLE VII. - THERMODYNAMIC PROPERTIES OF CESIUM ON SATURATION LINE

Ť,	P,	Partis	Partial pressure,	atm	x ₂ ,	Mg,		۷,	ρ,
м _о	atm	p ₁	ď	p2	g/g mixture	g/mole		g/oo	g/1000 cc
100 200 298-15	0.700238E-36 0.282767E-15 0.150230E-08	0.700238E-36 0.282767E-15 0.150223E-08		0 0.163606E-22 0.611635E-13	0 0.000000 0.000081	132,9100 132,9100 132,9154	0.43	0.881712E 38 0.436690E 18 0.122528E 12	0.113416=-34 0.2289958-14 0.8161428-08
a301.8 a301.8	C.219558E-08 C.219558E-08 C.219561E-08			0.104792E-12	0.000095	132,9163	000		0.117836F-07 0.117837E-07 0.117837E-07
004	0.295484E-C3 0.295484E-C3 0.557482E-C2	0.341672E-05 0.294083E-03 0.549759E-02		0.296875E-08 0.140152E-05 0.772344E-04	0.001(35 0.009441 0.027330	133.5404 133.5404 134.7514	0.10	0.103981E 07 0.655416F 05	0.136390E-04 0.961717E-03 0.152575E-01
700 800	0.444664E-01 0.208062E-00			0.126485E-02 0.979553E-02	0.055316	136.6906	0.94	0.945055E 04 0.226720E 04	0.105814F-10 0.441072E-00
900 b937.71	0.683671E 00 0.100008E 01		88	0.463169E-01 0.757002E-01	0.140736	142.9706	0.53	0.538170E 03	
1000	0.175554E 01 0.3768C7E 01		0 0		0.162950	147.3621	0.16		
1200	0.707012E 01	0.617109E	6 6	.899025E 00	0.225626	149.8104 152.0058	0.02	0.929698E 02	0.107562E 02 0.179461E 02
1400 1500	.186618E		00		0.273202 0.291861	153.9381 155.6196	0.39	.399906E 02 .269923E 02	0.250059E 02 0.344919E 02
T, ^o K	h _C ,	Δh _{vap} ,	hg, cal/g	s _c , cal/(g)(^o K)	$\frac{\Delta s_{\text{vap}}}{\text{cal/(g)(}^{\text{O}}\text{K)}}$		^S g' cal/(g)(^O K)	$(c_p)_{\mathrm{eq}}$, $\mathrm{cal/(g)(^0K)}$	$(c_{ m p})_{ m fr},$ cal/(g)($^{ m O}$ K)
100 200 298-15	-9.7743 -4.9891 0	142.2668 141.2196 139.8963	132,4925 136,2305 139,8963	0.0968139 0.1298676 0.1501256	9 1.4226679 6 0.7060981 6 0.4652146		1.5194868 0.8359656 0.6193402	0.0373794 0.0373801 0.0375824	0.0373794
. 00	0.0971	139.8683	139.9653	0.1504492			0.6166767	0.0375969	0.0373792
a3c1.8	4.1043	135.9279	140.0321	0.1637288	8 0.4503905 8 0.3350579		0.6141193	0.0376113	0.0373791
2009	15.2216	131.8333	147.0549	0.1920465			0.4557130	0.0458544	
700	26.4397	126.1692	152.6089 154.8824	0.2109194			0.3911611	0.0614825	
900 b937.71	37 . 6578 39 . 7735	119.3871	157.0449	0.2250162	2 0.1326523 5 0.1259323		0.3576685	0.0676452	
	43.2676	115.9676	159.2352	0.2309262			0.3468938	0.0671774	0.0371375
1200	54.4857	109.5154	164.0011	0.2411527			0.3324155	0.0632733	0.0371519
1300	60.0948	106.5303	166.6251 169.4098	0.2456422	2 0.0319464 1 0.0740757		0.3275956	0.0609153 0.0586486	0.0371916
1500	71.3129	101.0281	172.3410	0.2536686	0		0.3210207	0.0565811	0.0373430

 $^{^{}m a}_{
m Melting}$ point. $^{
m b}_{
m Normal}$ boiling point to equilibrium mixture.

$$(\Delta H_{298.15}^{O})_{2} = (H_{298.15}^{O})_{Cs_{2}(g)} - 2(H_{298.15}^{O})_{Cs_{1}(g)}$$
 (2)

Equilibrium Constants

Two sets of logarithms of the equilibrium constants for the two formation reactions discussed previously are also listed in tables II to IV. The equilibrium constant K_f for formation from the assigned reference element (col. 8) is obtained from the standard free energy change (ΔF_T^O) by means of the equation

$$\log_{10} K_{f} = -\frac{(\Delta F_{T}^{0})}{2.3025851 \text{ RT}}$$
(3)

The equilibrium constant K for formation from the atomic gas is obtained from a similar equation

$$\log_{10} K = -\frac{\Delta F_{T}^{O}}{2.3025851 \text{ RT}}$$
 (4)

Atomic weights, the universal gas constant and the constants used in the evaluation of the entropy constant were the same as those used in reference 1.

SFIFCTION OF INITIAL DATA

Crystal

Heat capacities of the crystal (table I, p. 4) were derived by smoothing the experimental data from the sources given in table VIII.

The following procedure was used to obtain smooth data from 0^{O} to 100^{O} K. A least-squares fit of the data of reference 2 was used in the temperature range from 0^{O} to 4^{O} K. Data from 4^{O} to 20^{O} K were obtained from a curve drawn through the data of reference 2 at 4^{O} K and through the data of references 3 and 4 at 20^{O} K. Data from 20^{O} to 100^{O} K were obtained from a least-squares fit of the data of reference 4 in that temperature range.

Between 100° and 210° K, the data of reference 4 are approximately linear in tem-

TABLE VIII, - REFERENCES FOR HEAT

CAPACITIES OF CRYSTAL

Temperature range,	Source
0.3 to 4	Ref. 2
8 to 12	Private communication from
	Professor D. C. McCollum
	of University of California
	in Riverside, California
	and ref. 3
20 to 210	Ref. 4

perature. The heat capacities at 100^{O} and 210^{O} K yield the equation

$$C_p^0 = 5.76 + 0.004 \text{ T}$$
 (5)

Above 210° K, thermal effects associated with the cesium-oxygen system cause anomalous increases in heat capacity and therefore these data were not used. Equation (5) was used to extrapolate heat capacities up to the melting point. The melting point of 301.8° K and the

heat of fusion of 520.1 calories per mole were taken from reference 5.

Reference 4 gives a melting point of 300. 5° K and the total heat absorption between the crystal at 262. 5° K and the liquid at 300. 5° K. If this heat is compared with the sensible enthalpy change of the crystal between these two temperatures, as calculated from equation (5), the implied heat of melting is 514 calories per mole. This value is in reasonable agreement with the value adopted.

The enthalpy from the linear equation for C_p^O , when combined with the heat of fusion, gave

$$(H_{301.8}^{O})_{Cs(liquid)} - (H_{298.15}^{O})_{Cs(crystal)} = 545.5$$
 cal/mole (6)

Liquid

A number of empirical equations have been used to satisfactorily fit enthalpy data as a function of temperature for various substances; however, heat capacity data derived from differentiation of the enthalpy equations are often very unreliable. For any particular species, various empirical enthalpy equations that might be chosen yield considerably different sets of heat capacities.

Two sets of experimental enthalpy data were considered: that of references 6 and 7. The curve of reference 6 exhibits an anomalous bump from 351° to 620° K, followed by a straight line to 1238° K. It is recommended in reference 6 that just the straight-line portion be used, which has a standard deviation $\sigma = 1.5$ percent. For the cubic equation (from 340° to 1176° K) of reference 7, $\sigma = 1.1$ percent.

The implied $C_p^0 = 7.25$ calories per mole per 0K of reference 6 differs greatly from the heat capacities derived from the enthalpy curve of reference 7 and listed in table IX.

According to reference 7, the curves of heat capacity against temperature of liquid

TABLE IX. - HEAT CAPACITIES

DERIVED FROM ENTHALPY

CURVE OF REFERENCE 7

Absolute	Heat capacity at
temperature,	constant pressure for
T,	standard state,
°К	C ^o p,
	cal/(mole)(^O K)
301.8	10.35
500	7. 28
700	6. 23
1000	8, 53
1200	12.64
1500	22. 67

cesium and other alkali metals have a parabolic shape; that is, heat capacity initially decreases with increasing temperature, reaches a minimum, and then increases continuously. However, a parabolic shape for heat capacity of these species is only the result of selecting a cubic to represent experimental enthalpy data.

It was thought advisable to use a straight line for enthalpy because (1) the straight line, in the standard deviation sense, represents cesium enthalpies about as well as the cubic does and (2) the cubic has the disadvantage of implying very large heat capacities for even a few hundred degrees extrapolation.

The actual straight line used in this report was based on the straight-line data of reference 6 but was

constrained to go through the previously calculated relative heat content of 545.5 calories per mole at the melting point (eq. (6)). This gives the following equation:

$$(H_T^0)_{Cs(liquid)} - (H_{298. \ 15}^0)_{Cs(crystal)} = -1704.5 + 7.455 T$$
 (7)

Equation (7) implies a constant value of $C_p^0 = 7.455$ calories per mole per 0 K.

Had the data of reference 7 been fitted to a straight line and similarly constrained to go through the heat content of 545.5 calories per mole, a constant value of $C_p^O = 7.595$ calories per mole per O K would have resulted. The enthalpies would be, on the average, only about 1 percent higher than those derived from the data of reference 6.

The enthalpy data of reference 8 are about 6 to 7 percent lower than the data of both references 6 and 7 over the range from about 600^{0} to 1300^{0} K; therefore, they were not used.

Monomer (Gas)

The partition function used to compute the thermodynamic properties of Cs_1 was truncated by the temperature-dependent cutoff technique used in reference 1. The results of these computations are exactly the same up to 2200° K as would be obtained by simply using all levels given in reference 9 with no cutoff. Above 2200° K enthalpy and entropy differ only in the last place computed and heat capacity in the last or second last place.

Dimer (Gas)

The method used to compute the thermodynamic properties of Cs_2 is the method for diatomic molecules described in reference 1. The constants for the dimer are shown in table V (p. 8). They were taken from references 10 and 11.

Vapor Pressures

During the period 1913 to 1937, the investigations reported in references 12 to 18 obtained vapor pressure data for cesium in the temperature range from 238° to 670° K. This early work was evaluated in reference 19 where it was concluded that, except for the results of references 17 and 18, the early work is not too reliable.

The vapor pressures in the 455° to 589° K range (ref. 17) were used in reference 19 to generate a vapor pressure equation. Inasmuch as the data of reference 17 are only relative pressure data (obtained by the magneto-optical method) and are converted to absolute pressure data by relating them to the unreliable data of reference 13, it seems prudent to disregard the data of reference 17.

Recently, the first experimental measurements in the temperature range that included the boiling point $(742^{\circ} \text{ to } 1199^{\circ} \text{ K})$ were reported in reference 20. These data check reasonably well with the later data of reference 6 $(727^{\circ} \text{ to } 1334^{\circ} \text{ K})$.

In view of the previous observations, the data of references 18 (vapor pressures in the liquid range), 20, and 6 were combined and fitted by the least-squares technique. All the data were given the same weight. The results of several empirical equations were compared. The following equation, which is in the form of the Kirchhoff equation, was accepted because it gave as satisfactory results as the other forms considered:

$$\log_{10} P = -\frac{3920.38}{T} + 5.71342 - 0.519781 \log_{10} T$$
 (8)

If only vapor pressures were considered in this report, then the smoothed experimental vapor pressures obtained from equation (8) would be accepted. As pointed out in the INTRODUCTION, however, these smoothed data, when combined with other thermodynamic data, lead to inconsistencies in the heat of condensation of the monomer at 0^{O} K (ΔH_{O}^{O}). Therefore, the data from equation (8) are used only for preliminary analysis, whereas the final recommended vapor pressures are calculated as explained in detail in the section Pressures and Weight Fractions of Gaseous Species.

Heat of Dissociation at 00 K

An upper bound (4020 cm⁻¹) for the heat of dissociation at $0^{\rm O}$ K(D $_{\rm O}^{\rm O}$) is estimated in reference 21 by analytical extrapolation of the ground state vibrational intervals. A lower bound (3318 cm⁻¹) is estimated in reference 21 by making certain assumptions as to the final states of the excited atoms produced. The average of these two values, which is equivalent to 10 500±1000 calories per mole of dimer, is recommended in this reference.

The estimate of 0.453 eV or 10 450 calories per mole from the correlation of force constants, equilibrium internuclear distances, and heats of dissociation of the alkali metals is obtained in reference 22. A technique for deducing a value of D_{O}^{O} where D_{O}^{O} was not well known is suggested in reference 23. This technique gave an optimum D_{O}^{O} = 10 000 calories per mole which is in fairly close agreement with the spectroscopic value of 10 500 calories per mole.

The calculation of $(\Delta H_O^O)_C$, which is discussed in the next section, depends in part on the value of D_O^O . While no value of D_O^O that was tried gave a constant value of $(\Delta H_O^O)_C$ over the entire temperature range of 301.8° to 1500° K, values of D_O^O from about 10 000 to 11 500 calories per mole gave the most nearly constant value of $(\Delta H_O^O)_C$.

Based on this analysis, the conclusion was that the spectroscopic value was approximately correct, and the value $D_{\mathbf{O}}^{0} = 10\ 500$ calories per mole was selected.

Heat of Condensation of Monomer at 00 K

The values of $(\Delta H_O^0)_c$ in calories per mole are given in a number of publications as follows: -19 048 (ref. 24), -18 790 (ref. 19), -19 050 (ref. 25), and -19 035 (ref. 26).

The aforementioned values of $(\Delta H_O^O)_c$ are based on the pre-1962 vapor pressures. Therefore, this report will recommend a "best" (ΔH_O^O) of -18 920 on the basis of the newer vapor pressures, as discussed in the following section.

THERMODYNAMIC ANALYSIS

As discussed in the INTRODUCTION, the procedure for the selection of a value for $(\Delta H_O^O)_C$ is to find that $(\Delta H_O^O)_C$ for which the difference between the calculated and smoothed experimental vapor pressures is a minimum. As a first step, it is necessary to compute $(\Delta H_O^O)_C$ for each temperature in the range of interest.

Derivation of $(\Delta H_c^0)_c$

The zero-degree heat of condensation of monomer was computed by means of the relation

$$(\Delta H_{O}^{O})_{c} = (\Delta F_{T}^{O})_{c} - \Delta (F_{T}^{O} - H_{O}^{O})_{c}$$
 (9)

where

$$(\Delta F_T^0)_c = (F_T^0)_{Cs(c)} - (F_T^0)_{Cs_1(g)} = RT \ln p_1$$
 (10)

and

$$\Delta(F_{T}^{o} - H_{O}^{o})_{c} = (F_{T}^{o} - H_{O}^{o})_{Cs(c)} - (F_{T}^{o} - H_{O}^{o})_{Cs_{1}(g)}$$
(11)

Partial Pressure of Monomer

The partial pressure of monomer $\,p_1^{}$ in atmospheres is obtained from the simultaneous solution of Dalton's Law and the equilibrium-constant equation for the dimerization reaction

$$2 \operatorname{Cs}_1 = \operatorname{Cs}_2$$

The equations are

$$P = p_1 + p_2 (12)$$

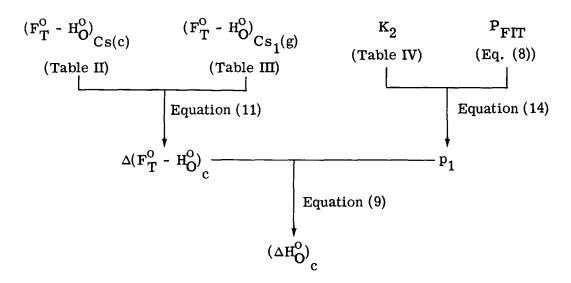
and

$$K_2 = \frac{p_2}{p_1^2} \tag{13}$$

where P is the total pressure, p_2 is the partial pressure of dimer, and K_2 is the equilibrium constant of dimerization. Values of K_2 are obtained from table IV (p. 7). The solution of equations (12) and (13) is

$$p_1 = \frac{-1 + \sqrt{4PK_2 + 1}}{2K_2} \tag{14}$$

Therefore, apparent values of $(\Delta H_O^O)_C$ can be computed at every temperature of interest by starting with given values of D_O^O , values of F_T^O - H_O^O for monomer, dimer, and condensed phase, and smoothed experimental vapor pressures P_{FIT} . This procedure can be summarized as follows:



RESULTS OF CALCULATING (ΔH_0^0)

The apparent $(\Delta H_O^0)_c$ values computed by means of this procedure from equation (9) are given in table X. The maximum difference is about 1 percent.

If some single arbitrary value of $(\Delta H_O^0)_c$ is selected, then for each of the previous temperatures there will be a difference between the smoothed experimental vapor pressure P_{FIT} and the vapor pressure P_{CALC} calculated with this arbitrary $(\Delta H_O^0)_c$ value.

For this analysis, an error function may be defined as $|P_{FIT} - P_{CALC}|/P_{FIT}$ averaged over the previous number of temperatures. The procedure for selecting a value of $(\Delta H_O^0)_c$ is to find that value of $(\Delta H_O^0)_c$ which gives the minimum error function.

The results of calculating the error function for various values of $(\Delta H_O^O)_c$ are given in table XI. The optimum $(\Delta H_O^O)_c$ was taken to be -18 920 calories per mole to the nearest 5 calories.

The fact that there is still about 4 percent error between the smoothed experimental vapor pressure and vapor pressure calculated with the optimum $(\Delta H_O^0)_c$, indicates that

Table X. - apparent values of $-(\Delta H_{\mathbf{O}}^{\mathbf{o}})_{\mathbf{c}}$ as a function of temperature

Tempera- ture, T, OK	$-(\Delta H_{O}^{O})_{c}$, cal/mole	Tempera- ture, T,	$-(\Delta H_{O}^{O})_{c}^{c}$, cal/mole
301.8	18 791. 6	1000	18 965.1
400	18 872.5	1100	18 954.3
500	18 921.7	1200	18 937.0
600	18 949.3	1300	18 913.1
700	18 963.6	1400	18 881.9
800	18 969.7	1500	18 843.4
900	18 970.0		

TABLE XI. - ERROR FUNCTION
FOR VARIOUS VALUES

OF
$$-(\Delta H_O^0)_c$$

$-(\Delta H_{O}^{O})_{c}^{C}$, cal/mole	Error function, ΔP /P
18 850	0.0538
18 900	. 0397
18 910	. 0381
18 920	. 0370
18 925	. 0370
18 930	. 0372
18 950	. 0389
18 970	. 0458

discrepancies still remain which may be in any of the following data:

- (a) Experimental vapor pressures
- (b) **D**O
- (c) Molecular constants of dimer (and, consequently, its thermodynamic properties)
- (d) Heat content (and other derived thermodynamic properties) of the condensed phase
- $(e) (\Delta H_{\mathbf{O}}^{\mathbf{O}})_{\alpha}$

The 4 percent error may also be due, in part, to the presence of a higher polymer, for example, trimer or tetramer; however, this possibility has not been considered in this report.

CALCULATION OF EQUILIBRIUM MIXTURE PROPERTIES ON SATURATION LINE

As explained in the INTRODUCTION, selection of an optimum value of $(\Delta H_O^0)_c$ makes it possible to compute a set of vapor pressures consistent with the sensible free energies of the pertinent species. These vapor pressures and other equilibrium properties on the saturation line (e.g., enthalpy, specific heat, and entropy) are given in table VII (p. 9). The equations used to obtain these properties are given in the following sections.

Pressures and Weight Fractions of Gaseous Species

By starting with the optimized $(\Delta H_O^0)_c$ and the sensible free energies of monomer and condensed phase, monomer pressure was computed with equations (9) to (11).

Dimer pressure was computed from equation (13). The total pressure P_{CALC} was obtained from the sum of p_1 and p_2 .

The weight fraction of monomer $\mathbf{x_1}$ was obtained from the mole fraction of monomer $\mathbf{N_1}$; thus,

$$N_1 = \frac{P_1}{P} \tag{15}$$

$$x_1 = \frac{N_1}{2 - N_1} \tag{16}$$

The weight fraction of the dimer x_2 was obtained from the defining equation

$$x_1 + x_2 = 1 (17)$$

Enthalpy of Vaporization and Entropy of Vaporization

The heat of vaporization in calories per gram from the condensed state to the equilibrium mixture on the saturation line was obtained from

$$\Delta h_{\text{vap}} = \frac{2x_1(\Delta H_T^0)_{f, 1} + x_2(\Delta H_T^0)_{f, 2}}{M_2}$$
 (18)

where M_2 is the molecular weight of dimer (265.82 g/mole). The heats of formation are obtained from tables III and IV (pp. 6 and 7).

The entropy of vaporization in calories per gram is obtained from equation (18) by means of

$$\Delta s_{\text{vap}} = \frac{\Delta h_{\text{vap}}}{T} \tag{19}$$

Enthalpy, Specific Heat, and Entropy of Cesium Vapor

The enthalpy h_g in calories per gram of an equilibrium mixture can be defined by

$$h_{g} = \sum_{i=1}^{2} \frac{x_{i}(H_{T}^{0})_{i}}{M_{i}} = \sum_{i=1}^{2} x_{i}h_{i}$$
 (20)

Equation (20) for saturation conditions gives the same results as the usual equation for the enthalpy of a mixture

$$h_g = h_{vap} + h_c$$

where h_c is the enthalpy of the condensed phase in calories per gram (table VII, p. 9). The equilibrium specific heat at constant pressure in calories per gram per ${}^{O}K$ is defined by

$$(c_p)_{eq} = \left(\frac{\partial h}{\partial T}\right)_p$$
 (21)

Using equation (20) results in equation (21) becoming

$$(c_{p})_{eq} = \sum_{i=1}^{2} x_{i}(c_{p})_{i} + \sum_{i=1}^{2} h_{i} \left(\frac{\partial x_{i}}{\partial T}\right)_{p}$$
(22)

or

$$(c_p)_{eq} = (c_p)_{fr} + (c_p)_r$$
(23)

From equation (17), it follows that

$$\left(\frac{\partial \mathbf{x_1}}{\partial \mathbf{T}}\right)_{\mathbf{p}} = -\left(\frac{\partial \mathbf{x_2}}{\partial \mathbf{T}}\right)_{\mathbf{p}} \tag{24}$$

Therefore,

$$(c_p)_r = \left(\frac{\partial x_2}{\partial T}\right)_p (h_2 - h_1) = \left(\frac{\partial x_2}{\partial T}\right)_p \frac{(\Delta H_T^0)_2}{M_2}$$
 (25)

Values of $(\Delta H_T^0)_2$ are tabulated in table IV (p. 7). Values of $(\partial x_2/\partial T)_p$ may be calculated from the following equation:

$$\left(\frac{\partial x_2}{\partial T}\right)_{p} = \frac{(\Delta H_T^0)_2}{RT^2} \frac{x_2(2 - x_2)(1 - x_2)}{2}$$
(26)

Equation (26) may be derived from the following expression, which is equivalent to equation (13):

$$K_2 = \frac{x_2(2 - x_2)}{4P(1 - x_2)^2} \tag{27}$$

and from the van't Hoff isobar, namely,

$$\left(\frac{\partial \ln K_2}{\partial T}\right)_p = \frac{(\Delta H_T^0)_2}{RT^2}$$
 (28)

The entropy of the gas mixture may be obtained from

$$s_g = \Delta s_{\text{vap}} + s_c \tag{29}$$

where s_c is the entropy of the condensed phase in calories per gram per ${}^{O}K$ (table VII, p. 9) and Δs_{vap} is obtained from equation (19). Equation (29) gives the identical results as the usual equation for the entropy of a mixture.

Molecular Weight of Mixture

An expression for the molecular weight of the mixture $\,M_g$ (in g/mole) in terms of the molecular weight and weight fraction of the dimer may be obtained as follows:

$$M_{g} = \frac{1}{n} = \frac{1}{n_{1} + n_{2}} = \frac{1}{\frac{x_{1}}{M_{1}} + \frac{x_{2}}{M_{2}}} = \frac{M_{2}}{2 - x_{2}}$$
(30)

where n, n_1 , and n_2 are the number of moles of mixture, monomer, and dimer per gram of mixture, respectively.

Specific Volume and Density of Mixture

The specific volume v in cubic centimeters per gram may be obtained from the ideal gas law

$$v = \frac{RT}{PM_g} = \frac{82.05971 \text{ T}}{PM_g}$$
 (31)

The density ρ in grams per 1000 cubic centimeters is, therefore,

$$\rho = \frac{1000}{v} \tag{32}$$

DISCUSSION OF VAPOR PRESSURES

Vapor Pressures Above Liquid

Table VI contains boiling points of liquid cesium at saturation pressures from 10⁻⁸ to 10 atmospheres at every power of 10. These data were calculated in two ways:

- (1) From the recommended thermodynamic data, as explained in the section THERMODYNAMIC ANALYSIS
- (2) By a least-squares equation fitted to the recommended P_{CALC} values with a standard deviation of 0.5 percent

$$\log_{10} P_{\text{atm}} = \frac{-4053.30}{T} + 7.04453 - 0.915282 \log_{10} T$$
 (33)

On the average the two sets of boiling points differ by about 0.2° K. It should be emphasized that the vapor pressure equation (eq. (33)) is recommended rather than equation (8).

Figures 1 and 2 compare the experimental vapor pressures of liquid cesium with a curve drawn through the recommended vapor pressures. The agreement of the two sets is seen to be very good. Figure 1 covers the temperature range from $303^{\rm O}$ to $1334^{\rm O}$ K. Figure 2 covers only the range from $729^{\rm O}$ to $1334^{\rm O}$ K to facilitate detailed comparison.

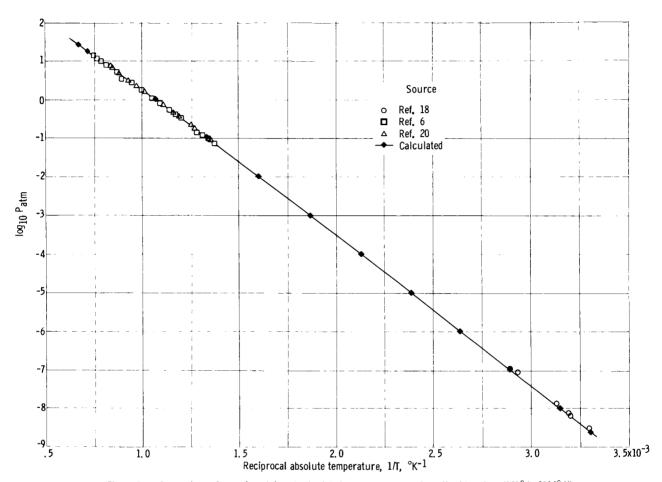


Figure 1. - Comparison of experimental and calculated vapor pressures above liquid cesium (303 $^{\circ}$ to 1334 $^{\circ}$ K).

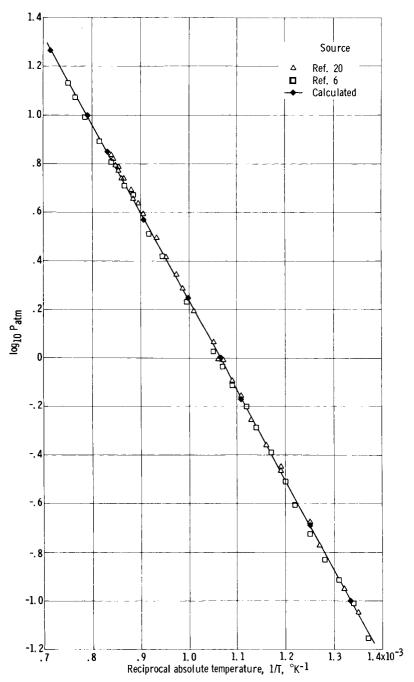


Figure 2. – Comparison of experimental and calculated vapor pressures above liquid cesium (729° to 1334° K).

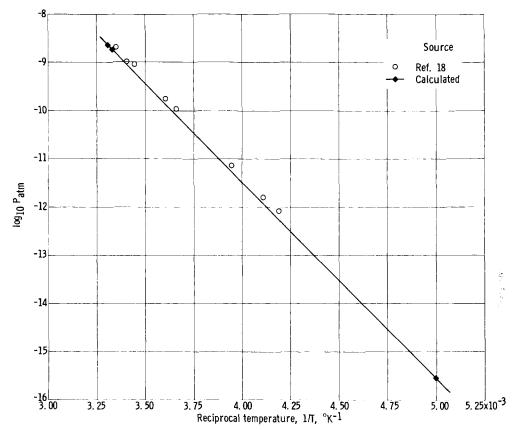


Figure 3. - Comparison of experimental and calculated vapor pressures above crystalline cesium.

Subsequent to the completion of the analysis, a vapor pressure equation in the temperature range 928° to 1558° K was presented in reference 27. On the average, these vapor pressures are about 2 percent lower than those recommended in this report.

Vapor Pressures Above Crystal

Experimental vapor pressures above the crystal are available from reference 18 at temperatures from 238° to 298° K. In order to be consistent with the thermodynamic data adopted for this report, however, vapor pressures from 100° to 301.8° K were calculated in the same manner as were the vapor pressures above the liquid.

On the average, the calculated vapor pressures between $238^{\rm O}$ and $298^{\rm O}$ K were 30 percent lower than the experimental ones (fig. 3). This difference is not significant, however, inasmuch as at the low pressure levels involved, the average difference is only in the order of 10^{-9} atmosphere. Furthermore, the lower calculated vapor pres-

sures are consistent with the assumption of reference 26 that the experimental vapor pressures are too high.

Lewis Research Center,
National Aeronautics and Space Administration,
Cleveland, Ohio, April 26, 1965.

APPENDIX - SYMBOLS

Ве	rotational constant, cm ⁻¹	$(\Delta H_{\mathrm{T}}^{\mathrm{O}})_{\mathrm{f}}$	enthalpy change for formation
C _p	heat capacity at constant pressure for standard		of substance from assigned reference element, cal/mole
	state, cal/(mole)(^O K)	H_{O}^{O}	chemical energy at 0°K for standard state, cal/mole
(c _p) _{eq}	equilibrium specific heat at constant pressure of mix-	h	enthalpy, cal/g
(c _p) _{fr}	ture, cal/g frozen specific heat at con-	$^{\Delta h}$ vap	heat of vaporization to 1 gram of equilibrium vapor
	stant pressure of mixture, cal/g	K	equilibrium constant for re- action of formation from
$(c_p)_r$	reactive specific heat at con- stant pressure of mixture,	K _f	element in atomic gas state equilibrium constant for re-
D	cal/g	1	action of formation from assigned reference element molecular weight based on chemical scale of natural oxygen, g/mole
D _e	spectroscopic constant for rotational stretching, cm ⁻¹	M	
$D_{O}^{\mathbf{O}}$	dissociation energy at 0° K, cal/mole		
$\mathbf{F_{T}^{o}}$	Gibbs free energy for stand- ard state, cal/mole	N _i	mole fraction of i th species, moles of i th species per
F_{T}^{O} - H_{O}^{O}	sensible free energy for standard state, cal/mole	n	mole of mixture number of moles per gram of equilibrium mixture
$H_{\mathbf{T}}^{\mathrm{O}}$	sum of sensible enthalpy at T ^O K and chemical energy	P	total vapor pressure, atm
	at 0° K for standard state, cal/mole	$p_{\dot{1}}$	partial pressure of i th species, atm
$H_T^0 - H_O^0$	sensible enthalpy for stand- ard state, cal/mole	R	universal gas constant, 1. 98726 cal/(mole)(^O K) or 82. 05971 (cc)(atm)/ (mole)(^O K)
$\Delta H_{\mathrm{T}}^{\mathrm{o}}$	enthalpy change for formation of substance from element in atomic gas state, cal/mole		
		s_T^o	entropy for standard state, cal/(mole)(^O K)

s	entropy, cal/(g)(^O K)	$^{\omega}{}_{ m e}$	zero-order vibrational fre-
$\Delta s_{ ext{vap}}$	entropy of vaporization to 1 gram of equilibrium vapor		quency for diatomic molecule, cm ⁻¹
Т	absolute temperature, ^O K	$\omega_{\mathrm{e}}^{\mathrm{x}}_{\mathrm{e}}, \omega_{\mathrm{e}}^{\mathrm{y}}_{\mathrm{e}}$	anharmonicity constants for diatomic molecule, cm ⁻¹
v v	specific volume of vapor, cc/g weight fraction of i th species,	Subscripts:	,
^x i	g of i th species per g of	c	condensed phase property
	mixture	f	formation from assigned
$^{lpha}\mathrm{e}$	vibration-rotation interaction		reference element
	constant for diatomic mole- cule, cm ⁻¹	g	property of equilibrium vapor mixture
ρ	density of vapor, g/1000 cc	1	monomer property
σ	standard deviation	2	dimer property

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